

Gradient estimates for Dirac-Harmonic maps with curvature term

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Abstract. In this paper, we derive a gradient estimate and Liouville theorem for Dirac-harmonic maps with curvature term.

1. INTRODUCTION

Let (M, g) be a Riemannian spin manifold, denote ΣM the spin bundle on M . Let (N, h) be another Riemannian manifold, $\Phi : M \rightarrow N$ a map. On the twisted bundle $\Sigma M \otimes \Phi^{-1}TN$, there is a Hermitian product $\langle \cdot, \cdot \rangle$ and a connection ∇ induced from those on ΣM and $\Phi^{-1}TN$. Denote $\not{D}\Psi := e_i \cdot \nabla_{e_i} \Psi$ the *Dirac operator along the map* Φ , $\forall \Psi \in \Gamma(\Sigma M \otimes \Phi^{-1}TN)$, where $\{e_i, i = 1, 2, \dots, m\}$ is a local orthonormal frame on M . *Dirac-harmonic maps* ([2, 3]) are critical points (Φ, Ψ) of the following functional:

$$(1.1) \quad L(\Phi, \Psi) = \frac{1}{2} \int_M [\|d\Phi\|^2 + \langle \Psi, \not{D}\Psi \rangle].$$

This functional comes from the nonlinear supersymmetric sigma model of quantum field theory (see [9]), the only difference is that in physics, the components of the field Ψ take values in the Grassmannian algebra of infinitely dimension while in (1.1) the components of Ψ are just the usual spinors in spin geometry, leaving the model in the category of geometric variational problems.

The Euler-Lagrange equations for L are given by

$$(1.2) \quad \begin{aligned} \tau(\Phi) &= \frac{1}{2} \langle \Psi^\alpha, e_i \cdot \Psi^\beta \rangle R^N(\theta_\alpha, \theta_\beta) \Phi_*(e_i), \\ \not{D}\Psi &= 0, \end{aligned}$$

where $R^N(X, Y) := [\nabla_X^N, \nabla_Y^N] - \nabla_{[X, Y]}^N$, $\forall X, Y \in \Gamma(TN)$ is the curvature operator of N , $\{\theta_\alpha\}$ a local frame of TN , $\Psi = \Psi^\alpha \otimes \theta_\alpha$ and $\tau(\Phi) := (\nabla_{e_i}^{T^*M \otimes \Phi^{-1}TN} d\Phi)(e_i)$ is the tension field of Φ .

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In [5], the authors derived Liouville theorems for Dirac-harmonic maps, namely, all solutions (Φ, Ψ) of (1.2) from the Euclidean space \mathbb{R}^n ($n \geq 3$), the hyperbolic space \mathbb{H}^n ($n \geq 3$) and the Riemannian Schwarzschild space \mathfrak{S}^n ($n \geq 3$) to any Riemannian manifolds with finite energy must be $(Constant, 0)$.

In fact, the supersymmetric sigma model in superstring theory includes an additional curvature term in (1.1) (see [9]). In [5], the authors introduced the following functional:

$$(1.3) \quad L_c(\Phi, \Psi) = \frac{1}{2} \int_M [\|d\Phi\|^2 + \langle \Psi, \not{D}\Psi \rangle - \frac{1}{6} R_{\alpha\gamma\beta\zeta} \langle \Psi^\alpha, \Psi^\beta \rangle \langle \Psi^\gamma, \Psi^\zeta \rangle] .$$

Critical points (Φ, Ψ) of L_c are called *Dirac harmonic maps with curvature term*.

The Euler-Lagrange equations for the action functional L_c are given by

$$(1.4) \quad \begin{aligned} \tau(\Phi) &= \frac{1}{2} \langle \Psi^\alpha, e_i \cdot \Psi^\beta \rangle R^N(\theta_\alpha, \theta_\beta) \Phi_*(e_i) \\ &\quad - \frac{1}{12} \Phi^*(\nabla^{TN} R^N)_{\alpha\beta\gamma\zeta} \langle \Psi^\alpha, \Psi^\gamma \rangle \langle \Psi^\beta, \Psi^\zeta \rangle , \\ \not{D}\Psi &= \frac{1}{3} R(\Psi, \Psi) \Psi . \end{aligned}$$

Due to the new nonlinear terms in the equations, (1.4) is more subtle than (1.2). A Liouville theorem for these kind of maps was also established in [5].

On the other hand, gradient estimates for harmonic functions was developed by S.T. Yau in his seminal paper [14], later, it was extended to the case of harmonic maps by Cheng [6] and Choi [8]. The method of gradient estimates has become a fundamental tool for the existence and uniqueness problems of differential equations on manifolds. In [4], extending Cheng and Choi's results, the authors obtained a gradient estimate and Liouville theorem for Dirac-harmonic maps. Later, Branding [1] proved a gradient estimate for the nonlinear supersymmetric sigma model.

The aim of the present paper is to extending results in [4] to the case of Dirac-harmonic maps with curvature term, we will prove a gradient estimate (Theorem 2.1) with concise and clear dependence on the curvatures of both domain and target manifolds, which enables us to derive a Liouville theorem (Theorem 3.1) as an application.

2. GRADIENT ESTIMATES

In the sequel, we assume that (M, g) is a complete non-compact Riemannian spin manifold. In this section, we derive a gradient estimate for Dirac-harmonic maps with curvature term from M into a regular ball on arbitrary Riemannian manifold (N, h) .

2.1. Preliminaries. We first recall the following Weitzenböck formula for a map Φ (see e.g. [11]).

Proposition 2.1. *For a smooth map $\Phi : M \rightarrow N$,*

$$(2.1) \quad \begin{aligned} \frac{1}{2} \Delta \|d\Phi\|^2 &= \|\nabla d\Phi\|^2 + \langle \nabla_{e_i} \tau(\Phi), \Phi_*(e_i) \rangle + \langle \Phi_*(Ric^M(e_i)), \Phi_*(e_i) \rangle \\ &\quad - R^N(\Phi_*(e_i), \Phi_*(e_j), \Phi_*(e_i), \Phi_*(e_j)) , \end{aligned}$$

where $\{e_i\}$ is a local orthonormal frame on M .

For a harmonic spinor field along a map Φ , we also have the similar Weitzenböck formula (cf. [3], Proposition 3.4).

Proposition 2.2. *Let M and N be Riemannian manifolds, $\Phi : M \rightarrow N$, and $\Psi = \Psi^\alpha \otimes \theta_\alpha \in \Gamma(\Sigma M \otimes \Phi^{-1}TN)$. Then*

$$(2.2) \quad \begin{aligned} \frac{1}{2}\Delta\|\Psi\|^2 &= \|\nabla\Psi\|^2 - \langle \not{D}^2\Psi, \Psi \rangle + \frac{1}{4}S_M\|\Psi\|^2 \\ &\quad - \frac{1}{2}R_{\alpha\beta\gamma\zeta}\langle e_i \cdot \Psi^\alpha, e_j \cdot \Psi^\beta \rangle \phi_i^\gamma \phi_j^\zeta, \end{aligned}$$

here $\{e_i\}$ and $\{\theta_\alpha\}$ are the local orthonormal frame on M, N , respectively. S_M is the scalar curvature of M .

Now we recall the following Kato-Yau inequalities in [4], which generalize both the result for harmonic maps in [13] and the result for harmonic spinors in [10].

Proposition 2.3 (Kato-Yau inequality). *Let E be any Dirac bundle on M with dimension m . Then for any cross-section $\Psi \in \Gamma(E)$ and $\delta > 0$, we have*

$$(2.3) \quad \|\nabla\Psi\|^2 \geq \left(1 + \frac{1}{m-1+\delta}\right) \|\nabla\|\Psi\|^2 - \frac{1}{\delta}\|\not{D}\Psi\|^2,$$

provided that $\Psi \neq 0$.

In particular, for any map $\Phi : M \rightarrow N$ and $\delta > 0$, we have

$$(2.4) \quad \|\nabla d\Phi\|^2 \geq \left(1 + \frac{1}{m-1+\delta}\right) \|\nabla\|d\Phi\|^2 - \frac{1}{\delta}\|\tau(\Phi)\|^2.$$

Lemma 2.1. *Assume that the pair (Φ, Ψ) is a smooth solution of (1.4). Suppose that the Ricci curvature of M satisfies $\text{Ric}_M \geq -\kappa$ for some nonnegative constant κ , the sectional curvature sec_N and the curvature tensor R^N of N satisfy $-b_2 \leq \text{sec}_N \leq b_1$, $\|\nabla R^N\| \leq b_3$ and $\|\nabla^2 R^N\| \leq b_4$ respectively, where b_i are constants with $b_2 \geq b_1 > 0, b_3 \geq 0, b_4 \geq 0$. Then*

$$(2.5) \quad \begin{aligned} |\langle \nabla_{e_i}\tau(\Phi), \Phi_*(e_i) \rangle| &\leq b_3 C(m, n)\|\Psi\|^2\|d\Phi\|^3 \\ &+ b_2 C(m, n)\|\Psi\|^2\|d\Phi\|\|\nabla d\Phi\| + b_2 C(m, n)\|\Psi\|\|\nabla\Psi\|\|d\Phi\|^2 \\ &+ b_4 C(m, n)\|\Psi\|^4\|d\Phi\|^2 + b_3 C(m, n)\|\Psi\|^3\|\nabla\Psi\|\|d\Phi\| \end{aligned}$$

for some constant $C(m, n) > 0$ depending only on m and n .

Proof. From the first equation in (1.4), we have

$$\begin{aligned} \langle \nabla_{e_i}\tau(\Phi), \Phi_*(e_i) \rangle &= \frac{1}{2} \langle \langle \Psi^\alpha, e_k \cdot \Psi^\beta \rangle (\nabla_{\Phi_*(e_i)} R^N)(\theta_\alpha, \theta_\beta) \Phi_*(e_k), \Phi_*(e_i) \rangle \\ &+ \frac{1}{2} \langle \langle \Psi^\alpha, e_k \cdot \Psi^\beta \rangle R^N(\theta_\alpha, \theta_\beta)(\nabla_{e_i} d\Phi)(e_k), \Phi_*(e_i) \rangle \\ &+ \frac{1}{2} \langle \langle \nabla_{e_i} \Psi^\alpha, e_k \cdot \Psi^\beta \rangle R^N(\theta_\alpha, \theta_\beta) \Phi_*(e_k), \Phi_*(e_i) \rangle \\ &+ \frac{1}{2} \langle \langle \Psi^\alpha, \nabla_{e_i}(e_k \cdot \Psi^\beta) \rangle R^N(\theta_\alpha, \theta_\beta) \Phi_*(e_k), \Phi_*(e_i) \rangle \\ &- \frac{1}{12} \langle (\nabla_{\Phi_*(e_p)} (\nabla R^N)_{\alpha\beta\gamma\zeta}) \langle \Psi^\alpha, \Psi^\gamma \rangle \langle \Psi^\beta, \Psi^\zeta \rangle, \Phi_*(e_p) \rangle \\ &- \frac{1}{3} \langle (\nabla R^N)_{\alpha\beta\gamma\zeta} \langle \Psi^\alpha, \Psi^\gamma \rangle \langle \Psi^\beta, \nabla_{e_p}(\Psi^\zeta) \rangle, \Phi_*(e_p) \rangle, \end{aligned}$$

thus (2.5) follows. □

Lemma 2.2. *If (Φ, Ψ) is a smooth solutions of (1.4), then*

$$(2.6) \quad \|\mathcal{D}^2\Psi\| \leq b_3 C(m, n) \|\Psi\|^3 \|d\Phi\| + b_2 C(m, n) \|\Psi\|^2 \|\nabla\Psi\| ,$$

for some constant $C(m, n) > 0$ depending only on m and n .

Proof. By the second equation in (1.4), it is easy to see that

$$\|\mathcal{D}^2\Psi\| \leq C(m, n) (\|\nabla R\| \|\Psi\|^3 \|d\Phi\| + \|R^N\| \|\Psi\|^2 \|\nabla\Psi\|) ,$$

so (2.6) holds. □

2.2. Gradient estimates. Combined with (2.5) and Weitzenböck formula (2.1), one can get that

$$\begin{aligned} \frac{1}{2}\Delta\|d\Phi\|^2 &\geq \|\nabla d\Phi\|^2 - b_3 C(m, n) \|\Psi\|^2 \|d\Phi\|^3 - b_2 C(m, n) \|\Psi\|^2 \|d\Phi\| \|\nabla d\Phi\| \\ &\quad - b_2 C(m, n) \|\Psi\| \|\nabla\Psi\| \|d\Phi\|^2 - b_4 C(m, n) \|\Psi\|^4 \|d\Phi\|^2 \\ &\quad - b_3 C(m, n) \|\Psi\|^3 \|\nabla\Psi\| \|d\Phi\| - \kappa \|d\Phi\|^2 - (1 - \frac{1}{p}) b_1 \|d\Phi\|^4 . \end{aligned}$$

Applying Cauchy-Schwarz inequality, it follows that, for any constant $\delta_1 > 0$ depending only on m, n ,

$$\begin{aligned} \frac{1}{2}\Delta\|d\Phi\|^2 &\geq (1 - \delta_1) \|\nabla d\Phi\|^2 - \frac{b_2^2}{\delta_1} C(m, n) \|\Psi\|^4 \|d\Phi\|^2 \\ &\quad - \frac{b_3^2}{b_1} C(m, n) 2p(p+1) \|\Psi\|^4 \|d\Phi\|^2 \\ &\quad - \frac{b_1}{2p(p+1)} \|d\Phi\|^4 - \frac{b_2^2}{b_1} C(m, n) 2p(p+1) \|\Psi\|^2 \|\nabla\Psi\|^2 \\ &\quad - \frac{b_1}{2p(p+1)} \|d\Phi\|^4 \\ &\quad - b_4 C(m, n) \|\Psi\|^4 \|d\Phi\|^2 - \frac{b_3^2}{b_1} C(m, n) \|\Psi\|^4 \|d\Phi\|^2 \\ &\quad - b_1 C(m, n) \|\Psi\|^2 \|\nabla\Psi\|^2 \\ &\quad - \kappa \|d\Phi\|^2 - (1 - \frac{1}{p}) b_1 \|d\Phi\|^4 . \end{aligned}$$

Since $b_2 \geq b_1 > 0$ and δ_1 depending only on m, n , we obtain

$$(2.7) \quad \begin{aligned} \frac{1}{2}\Delta\|d\Phi\|^2 &\geq (1 - \delta_1) \|\nabla d\Phi\|^2 - (b_2^2 + \frac{b_3^2}{b_1} + b_4) C(m, n) \|\Psi\|^4 \|d\Phi\|^2 \\ &\quad - \frac{b_2^2}{b_1} C(m, n) \|\Psi\|^2 \|\nabla\Psi\|^2 - \kappa \|d\Phi\|^2 - (1 - \frac{1}{p+1}) b_1 \|d\Phi\|^4 . \end{aligned}$$

By using (2.6) and Weitzenböck formula (2.2), we estimate $\frac{1}{2}\|\Psi\|^4$ as follows:

$$\begin{aligned}
\frac{1}{2}\Delta\|\Psi\|^4 &= \|\Psi\|^2\Delta\|\Psi\|^2 + \|\nabla\|\Psi\|^2\|^2 \\
&\geq 2\|\Psi\|^2\|\nabla\Psi\|^2 + \|\nabla\|\Psi\|^2\|^2 + \frac{S_M}{2}\|\Psi\|^4 - b_2^2 C(m, n)\|\Psi\|^4\|d\Phi\|^2 \\
&\quad - 2\langle \mathcal{D}^2\Psi, \Psi \rangle\|\Psi\|^2 \\
&\geq 2\|\Psi\|^2\|\nabla\Psi\|^2 + \|\nabla\|\Psi\|^2\|^2 - \frac{m}{2}\kappa\|\Psi\|^4 - b_2^2 C(m, n)\|\Psi\|^4\|d\Phi\|^2 \\
&\quad - b_3 C(m, n)\|\Psi\|^6\|d\Phi\| - b_2 C(m, n)\|\Psi\|^5\|\nabla\Psi\|.
\end{aligned}$$

Applying Cauchy-Schwarz inequality,

$$\begin{aligned}
\frac{1}{2}\Delta\|\Psi\|^4 &\geq 2\|\Psi\|^2\|\nabla\Psi\|^2 + \|\nabla\|\Psi\|^2\|^2 - \frac{m}{2}\kappa\|\Psi\|^4 \\
&\quad - b_2^2 C(m, n)\|\Psi\|^4\|d\Phi\|^2 \\
&\quad - b_3 C(m, n)\|\Psi\|^4\|d\Phi\|^2 - b_3 C(m, n)\|\Psi\|^8 \\
(2.8) \quad &\quad - \frac{1}{2}\|\Psi\|^2\|\nabla\Psi\|^2 - b_2^2 C(m, n)\|\Psi\|^8 \\
&= \frac{3}{2}\|\Psi\|^2\|\nabla\Psi\|^2 + \|\nabla\|\Psi\|^2\|^2 - \frac{m}{2}\kappa\|\Psi\|^4 \\
&\quad - (b_2^2 + b_3)C(m, n)\|\Psi\|^4\|d\Phi\|^2 - (b_3 + b_2^2)C(m, n)\|\Psi\|^8.
\end{aligned}$$

Putting Kato-Yau inequality (2.4) into (2.7), when $\delta = 1$. Set $\delta_1 = \frac{3}{4 + 4m}$ such that $(1 - \delta_1)(1 + \frac{1}{m}) = 1 + \frac{1}{4m}$, then

$$\begin{aligned}
\frac{1}{2}\Delta\|d\Phi\|^2 &\geq (1 + \frac{1}{4m})\|\nabla\|d\Phi\|\|^2 - (b_2^2 + \frac{b_3^2}{b_1} + b_4)C(m, n)\|\Psi\|^4\|d\Phi\|^2 \\
(2.9) \quad &\quad - b_3^2 C(m, n)\|\Psi\|^8 - \frac{b_2^2}{b_1} C(m, n)\|\Psi\|^2\|\nabla\Psi\|^2 \\
&\quad - \kappa\|d\Phi\|^2 - (1 - \frac{1}{p+1})b_1\|d\Phi\|^4.
\end{aligned}$$

As for $\frac{1}{2}\Delta\|\Psi\|^4$, because of $\|\nabla\Psi\|^2 \geq \frac{1}{m}\|\nabla\|\Psi\|\|^2$,

$$\begin{aligned}
\frac{1}{2}\Delta\|\Psi\|^4 &\geq (1 + \frac{1}{4m})\|\nabla\|\Psi\|^2\|^2 + \frac{1}{2}\|\Psi\|^2\|\nabla\Psi\|^2 - \frac{m}{2}\kappa\|\Psi\|^4 \\
(2.10) \quad &\quad - (b_2^2 + b_3)C(m, n)\|\Psi\|^4\|d\Phi\|^2 - (b_3 + b_2^2)C(m, n)\|\Psi\|^8.
\end{aligned}$$

Now we fix the constant $C_0 = C(m, n)$ in the above inequalities (2.9) (2.10). Set

$$C_1 = \sqrt{\frac{2b_2^2}{b_1} C_0} \quad , \quad e_1 = C_1\|\Psi\|^2 \quad , \quad e := \sqrt{\|d\Phi\|^2 + e_1^2} = \sqrt{\|d\Phi\|^2 + C_1^2\|\Psi\|^4} \quad ,$$

then

$$\begin{aligned}
\frac{1}{2}\Delta e^2 &= \frac{1}{2}\Delta\|d\Phi\|^2 + \frac{1}{2}\Delta C_1^2\|\Psi\|^4 \\
&\geq \left(1 + \frac{1}{4m}\right)\|\nabla\|d\Phi\|\|^2 + \left(1 + \frac{1}{4m}\right)C_1^2\|\nabla\|\Psi\|^2\|^2 \\
&\quad - \left(b_3^2 + \frac{b_2^2 b_3}{b_1} + \frac{b_2^4}{b_1}\right)C(m, n)\|\Psi\|^8 \\
&\quad - \left(b_2^2 + \frac{b_3^2}{b_1} + b_4 + \frac{b_2^4}{b_1} + \frac{b_2^2 b_3}{b_1}\right)C(m, n)\|\Psi\|^4\|d\Phi\|^2 \\
&\quad - \kappa\|d\Phi\|^2 - \frac{m}{2}\kappa\|\Psi\|^4 - \left(1 - \frac{1}{p+1}\right)b_1\|d\Phi\|^4 \\
&\geq \left(1 + \frac{1}{4m}\right)(\|\nabla\|d\Phi\|\|^2 + C_1^2\|\nabla\|\Psi\|^2\|^2) \\
&\quad - \left(b_2^2 + \frac{b_3^2}{b_1} + b_4 + \frac{b_2^4}{b_1} + \frac{b_2^2 b_3}{b_1} + b_3^2\right)C(m, n)\|\Psi\|^4 e^2 \\
&\quad - \left(1 + \frac{m}{2}\right)\kappa e^2 - \left(1 - \frac{1}{p+1}\right)b_1\|d\Phi\|^2 e^2 .
\end{aligned}$$

Since $\|\nabla e\|^2 \leq \|\nabla\|d\Phi\|\|^2 + \|\nabla e_1\|^2$, we have

$$\begin{aligned}
\Delta e &= e^{-1}\left(\frac{1}{2}\Delta e^2 - \|\nabla e\|^2\right) \\
(2.11) \quad &\geq \frac{1}{4m}\frac{\|\nabla e\|^2}{e} - \left(b_2^2 + \frac{b_3^2}{b_1} + b_4 + \frac{b_2^4}{b_1} + \frac{b_2^2 b_3}{b_1} + b_3^2\right)C(m, n)\|\Psi\|^4 e \\
&\quad - \left(1 + \frac{m}{2}\right)\kappa e - \left(1 - \frac{1}{p+1}\right)b_1\|d\Phi\|^2 e .
\end{aligned}$$

Let $r(x)$ be the distance function from the point x_0 in M . Define a function on the geodesic ball $B_a(x_0)$ by

$$F = (a^2 - r^2)f = (a^2 - r^2)\frac{e}{B \circ \Phi} ,$$

where $B : N \rightarrow \mathbb{R}^+$ is a function which we will be defined later. Clearly, the function F vanishes on the boundary $B_a(x_0)$, then F must achieve its maximum at an interior point x^* when $e \neq 0$. In addition, the distance function r is a smooth near the point x^* , therefore we may assume that the distance function r is twice differentiable near the point x^* (cf. [7]). By the maximum principle, we have

$$(2.12) \quad \nabla F(x^*) = 0 ,$$

$$(2.13) \quad \Delta F(x^*) \leq 0 .$$

Now we recall the Laplacian Comparison Theorem [12], for some constant $C(m) >$ depending only on m ,

$$(2.14) \quad \Delta r^2 \leq C(m)(1 + \sqrt{\kappa}r) .$$

Lemma 2.3. *At the point x^* , we have the following estimate*

$$\begin{aligned}
(2.15) \quad & \frac{1}{4m} \frac{\|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2} - \frac{\Delta(B \circ \Phi)}{B \circ \Phi} \\
& - 4r \frac{\|\nabla(B \circ \Phi)\|}{(a^2 - r^2)(B \circ \Phi)} - \tilde{b} C(m, n) \|\Psi\|^4 \\
& - \left(1 - \frac{1}{p+1}\right) b_1 \|d\Phi\|^2 - \left(1 + \frac{m}{2}\right) \kappa \\
& - \frac{C(m)(1 + \sqrt{\kappa}r)}{a^2 - r^2} - \frac{8r^2}{(a^2 - r^2)^2} \leq 0.
\end{aligned}$$

Here $\tilde{b} = b_2^2 + \frac{b_3^2}{b_1} + b_4 + \frac{b_2^4}{b_1} + \frac{b_2^2 b_3}{b_1} + b_3^2$.

Proof. We first have

$$(2.16) \quad \nabla f = \frac{\nabla e}{B \circ \Phi} - \frac{e \nabla(B \circ \Phi)}{(B \circ \Phi)^2},$$

and

$$(2.17) \quad \Delta f = \frac{\Delta e}{B \circ \Phi} - \frac{f \Delta(B \circ \Phi)}{B \circ \Phi} - \frac{2 \langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi}.$$

By using (2.11) and setting $\mu = \frac{1}{4m}$, $\tilde{b} = b_2^2 + \frac{b_3^2}{b_1} + b_4 + \frac{b_2^4}{b_1} + \frac{b_2^2 b_3}{b_1} + b_3^2$, we get

$$\begin{aligned}
(2.18) \quad \Delta f & \geq \mu \frac{\|\nabla e\|^2}{e(B \circ \Phi)} - \tilde{b} C(m, n) \|\Psi\|^4 f - \left(1 + \frac{m}{2}\right) \kappa f \\
& - \left(1 - \frac{1}{p+1}\right) b_1 \|d\Phi\|^2 f - \frac{f \Delta(B \circ \Phi)}{B \circ \Phi} - \frac{2 \langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi}.
\end{aligned}$$

Applying Cauchy-Schwarz inequality, we also have that

$$\begin{aligned}
(2.19) \quad -\frac{2 \langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} & = -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} \\
& - 2\mu \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} \\
& = -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} - 2\mu \frac{\langle \nabla(B \circ \Phi), \nabla e \rangle}{(B \circ \Phi)^2} \\
& + 2\mu \frac{f \|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2} \\
& \geq -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} - \mu \frac{\|\nabla e\|^2}{e(B \circ \Phi)} \\
& + \mu \frac{f \|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2}.
\end{aligned}$$

So we have

$$\begin{aligned}
\frac{\Delta f}{f} & \geq -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} + \mu \frac{\|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2} - \frac{\Delta(B \circ \Phi)}{B \circ \Phi} \\
& - \tilde{b} C(m, n) \|\Psi\|^4 - \left(1 + \frac{m}{2}\right) \kappa - \left(1 - \frac{1}{p+1}\right) b_1 \|d\Phi\|^2.
\end{aligned}$$

At the point x^* , since F achieves its maximum, as a consequence $\nabla F(x^*) = 0$ and $\Delta F(x^*) \leq 0$. By the definition of F , one has that

$$\frac{\nabla r^2}{a^2 - r^2} = \frac{\nabla f}{f},$$

and

$$-\frac{\Delta r^2}{a^2 - r^2} + \frac{\Delta f}{f} - \frac{2\langle \nabla r^2, \nabla f \rangle}{f(a^2 - r^2)} \leq 0.$$

It follows that

$$\frac{\Delta f}{f} - \frac{\Delta r^2}{a^2 - r^2} - \frac{2\|\nabla r^2\|^2}{(a^2 - r^2)^2} \leq 0.$$

Using $\|\nabla r^2\| = 2r$ and (2.14), we then have

$$0 \geq \frac{\Delta f}{f} - \frac{C(m)(1 + \sqrt{\kappa}r)}{a^2 - r^2} - \frac{8r^2}{(a^2 - r^2)^2},$$

i.e.

$$\begin{aligned} 0 &\geq -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{B \circ \Phi} + \mu \frac{\|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2} - \frac{\Delta(B \circ \Phi)}{B \circ \Phi} \\ &\quad - \tilde{b}C(m, n)\|\Psi\|^4 - (1 + \frac{m}{2})\kappa - (1 - \frac{1}{p+1})b_1\|d\Phi\|^2 \\ &\quad - \frac{C(m)(1 + \sqrt{\kappa}r)}{a^2 - r^2} - \frac{8r^2}{(a^2 - r^2)^2}. \end{aligned}$$

Since

$$\begin{aligned} -(2 - 2\mu) \frac{\langle \nabla(B \circ \Phi), \nabla f \rangle}{f(B \circ \Phi)} &= -(2 - 2\mu)2r \frac{\langle \nabla(B \circ \Phi), \nabla r \rangle}{(a^2 - r^2)(B \circ \Phi)} \\ &\geq -(2 - 2\mu)2r \frac{\|\nabla(B \circ \Phi)\|}{(a^2 - r^2)(B \circ \Phi)}, \end{aligned}$$

we conclude that

$$\begin{aligned} \mu \frac{\|\nabla(B \circ \Phi)\|^2}{(B \circ \Phi)^2} - \frac{\Delta(B \circ \Phi)}{B \circ \Phi} - 4r \frac{\|\nabla(B \circ \Phi)\|}{(a^2 - r^2)(B \circ \Phi)} - \tilde{b}C(m, n)\|\Psi\|^4 \\ - (1 - \frac{1}{p+1})b_1\|d\Phi\|^2 - (1 + \frac{m}{2})\kappa - \frac{C(m)(1 + \sqrt{\kappa}r)}{a^2 - r^2} - \frac{8r^2}{(a^2 - r^2)^2} \leq 0. \end{aligned}$$

□

Theorem 2.1 (Gradient Estimate). *Suppose the Ricci curvature of M satisfies $\text{Ric}_M \geq -\kappa$ for some nonnegative constant κ , the sectional curvature sec_N and curvature tensor R^N satisfy $-b_2 \leq \text{sec}_N \leq b_1$, $\|\nabla R^N\| \leq b_3$ and $\|\nabla^2 R^N\| \leq b_4$ respectively, where b_i are constants with $b_2 \geq b_1 > 0$, $b_3 \geq 0$, $b_4 \geq 0$. If (Φ, Ψ) is a solution of (1.4), $\Phi : M \rightarrow B_{y_0}(R) \subset N^n$, $R < \pi/(2\sqrt{b_1})$, then for any $x_0 \in M$ and any positive constant a , we have*

$$(2.20) \quad \sup_{B_{x_0}(\frac{a}{2})} \|d\Phi\| \leq \frac{C(m, n)}{\sqrt{b_1} \cos^2(\sqrt{b_1}R)} \left(\frac{1 + a\sqrt{\kappa}}{a} + \sqrt{\frac{b}{b_1}} \sup_{B_{x_0}(a)} \|\Psi\|^2 \right).$$

Here $b = b_1 b_2^2 + b_3^2 + b_4 b_1 + b_4^2 + b_2^2 b_3 + b_3^2 b_1 + b_1^{3/2} b_3$, and $C(m, n)$ is a constant depending only on the dimensions m and n .

Proof. Choose $B(y) = \sqrt{b_1} \cos(\sqrt{b_1} \rho(y))$ in the function $f = \frac{e}{B \circ \Phi}$, then

$$\begin{aligned} \Delta(B \circ \phi) &\leq (\text{Hess } B) \circ \Phi \|d\Phi\|^2 + \|(\nabla B) \circ \Phi\| \|\tau(\Phi)\| \\ &\leq -b_1^{3/2} \cos(\sqrt{b_1} \rho \circ \Phi) \|d\Phi\|^2 + b_1 C(m, n) (b_2 \|\Psi\|^2 \|d\Phi\| + b_3 \|\Psi\|^4) \\ &\leq -b_1^{3/2} \cos(\sqrt{b_1} \rho \circ \Phi) \|d\Phi\|^2 + \frac{1}{(p+1)(p+2)} b_1 B \circ \Phi \|d\Phi\|^2 \\ &\quad + \frac{1}{B \circ \Phi} (p+1)(p+2) b_1 b_2^2 C(m, n) \|\Psi\|^4 + b_1 b_3 C(m, n) \|\Psi\|^4. \end{aligned}$$

Putting this inequalities into (2.15), then we have

$$\begin{aligned} &\frac{b_1}{p+2} \|d\Phi\|^2 - \frac{4rb_1}{(a^2-r^2)B \circ \Phi} \|d\Phi\| \\ &\leq \tilde{b} C(m, n) \|\Psi\|^4 + \left(1 + \frac{m}{2}\right) \kappa \\ &\quad + \frac{C(m)(1+\sqrt{\kappa}r)}{a^2-r^2} + \frac{8r^2}{(a^2-r^2)^2} \\ &\quad + \frac{b_1 b_2^2}{(B \circ \Phi)^2} C(m, n) \|\Psi\|^4 + \frac{b_1 b_3}{B \circ \Phi} C(m, n) \|\Psi\|^4 \\ &\leq \left(1 + \frac{m}{2}\right) \kappa + \frac{C(m)(1+\sqrt{\kappa}r)}{a^2-r^2} + \frac{8r^2}{(a^2-r^2)^2} \\ &\quad + \frac{1}{(B \circ \Phi)^2} (b_1 b_2^2 + b_1^{3/2} b_3 + \tilde{b} b_1) C(m, n) \|\Psi\|^4. \end{aligned}$$

Namely,

$$\begin{aligned} \frac{b_1}{p+2} e^2 - \frac{4rb_1}{(a^2-r^2)B \circ \Phi} e &\leq \left(1 + \frac{m}{2}\right) \kappa + \frac{C(m)(1+\sqrt{\kappa}r)}{a^2-r^2} + \frac{8r^2}{(a^2-r^2)^2} \\ &\quad + \frac{1}{(B \circ \Phi)^2} (b_1 b_2^2 + b_1^{3/2} b_3 + \tilde{b} b_1) \\ &\quad \times C(m, n) \|\Psi\|^4. \end{aligned}$$

From the definition of F , we know that

$$\begin{aligned} &\frac{b_1(B \circ \Phi)^2}{p+2} F^2(x^*) - 4rb_1 F(x^*) \\ (2.21) \quad &\leq \left(\left(1 + \frac{m}{2}\right) \kappa + \frac{C(m)(1+\sqrt{\kappa}r)}{a^2-r^2} + \frac{8r^2}{(a^2-r^2)^2} \right. \\ &\quad \left. + \frac{1}{(B \circ \Phi)^2} (b_1 b_2^2 + b_1^{3/2} b_3 + \tilde{b} b_1) C(m, n) \|\Psi\|^4 \right) (a^2-r^2)^2. \end{aligned}$$

For the RHS of (2.21), denote $b = b_1 b_2^2 + b_1^{3/2} b_3 + \tilde{b} b_1$ and

$$\begin{aligned} I &:= \left(\left(1 + \frac{m}{2}\right) \kappa + \frac{C(m)(1+\sqrt{\kappa}r)}{a^2-r^2} + \frac{8r^2}{(a^2-r^2)^2} + \frac{1}{(B \circ \Phi)^2} b C(m, n) \|\Psi\|^4 \right) \\ &\quad \times (a^2-r^2)^2. \end{aligned}$$

We have the following estimate

$$\begin{aligned} I &\leq a^2 \left(\left(\left(1 + \frac{m}{2}\right)\kappa + \frac{1}{(B \circ \Phi)^2} C(m, n) b \|\Psi\|^4 \right) a^2 + C(m)(1 + \sqrt{\kappa} a) + 8 \right) \\ &\leq \frac{C(m, n) b_1}{(B \circ \Phi)^2} a^2 \left(1 + \sqrt{\kappa} a + \sqrt{\frac{b}{b_1}} \|\Psi\|^2 a \right)^2. \end{aligned}$$

It is elementary that if $Ax^2 - Bx - C \leq 0$ with A, B, C all positive, then

$$x \leq \frac{B}{A} + \sqrt{\frac{C}{A}}.$$

So we can conclude that

$$F(x^*) \leq \frac{C(m, n)a}{b_1 \cos^2(\sqrt{b_1} R)} \left(1 + \sqrt{\kappa} a + \sqrt{\frac{b}{b_1}} \sup_{B_{x_0}(a)} \|\Psi\|^2 a \right).$$

This proves Theorem 2.1. \square

3. LIOUVILLE THEOREM

Now using the gradient estimate, we can get a Liouville type theorem for Dirac-harmonic maps with curvature term.

Theorem 3.1 (Liouville Theorem). *Assume that M is complete with nonnegative Ricci curvature and the scalar curvature is bounded below by a positive constant ϵ , suppose the sectional curvature sec_N and curvature tensor R^N of N satisfy $-b_2 \leq \text{sec}_N \leq b_1$ and $\|\nabla R^N\| \leq b_3, \|\nabla^2 R^N\| \leq b_4$ respectively, where b_i are constants with $b_2 \geq b_1 > 0, b_3, b_4 \geq 0$. Then there is a constant $\delta > 0$ such that for any smooth solution (Φ, Ψ) of (1.4) satisfying $\Phi(M) \subset B_{y_0}(R) \subset N, R < \pi/(2\sqrt{b_1})$ and $\|\Psi\| < \delta$, we have $\Phi \equiv \text{constant}$ and $\Psi \equiv 0$.*

Proof. By using Weitzenböck formula (2.2) and (2.6), we can get that

$$\begin{aligned} \frac{1}{2} \Delta \|\Psi\|^2 &\geq \|\nabla \Psi\|^2 + \frac{\epsilon}{4} \|\Psi\|^2 - b_2^2 C(m, n) \|\Psi\|^2 \|d\Phi\|^2 - b_3 C(m, n) \|\Psi\|^4 \|d\Phi\| \\ &\quad - b_2 C(m, n) \|\Psi\|^3 \|\nabla \Psi\| \\ &\geq \|\nabla \Psi\|^2 + \frac{\epsilon}{4} \|\Psi\|^2 - b_2^2 C(m, n) \|\Psi\|^2 \|d\Phi\|^2 - b_3 C(m, n) \|\Psi\|^4 \|d\Phi\| \\ &\quad - \frac{1}{2} \|\nabla \Psi\|^2 - b_2^2 C(m, n) \|\Psi\|^6 \\ &\geq \left(\frac{\epsilon}{4} - b_2^2 C(m, n) \|d\Phi\|^2 - b_3 C(m, n) \|\Psi\|^2 \|d\Phi\| \right. \\ &\quad \left. - b_2^2 C(m, n) \|\Psi\|^4 \right) \|\Psi\|^2. \end{aligned}$$

Here $C(m, n)$ is a positive constant depending only on m, n . It is obvious that this theorem is valid when M is compact. Now we assume that M is non-compact. Suppose $\|\Psi\| < \delta$, then according to Theorem 2.1,

$$\|d\Phi\| \leq \frac{C(m, n) \sqrt{b} \delta^2}{b_1 \cos^2(\sqrt{b_1} R)}.$$

Therefore,

$$\frac{1}{2}\Delta\|\Psi\|^2 \geq \left(\frac{\epsilon}{4} - \frac{bb_2^2C(m,n)}{b_1^2\cos^4(\sqrt{b_1}R)}\delta^4 - \frac{\sqrt{b}b_3C(m,n)}{b_1\cos^2(\sqrt{b_1}R)}\delta^4 - b_2^2C(m,n)\delta^4 \right) \|\Psi\|^2.$$

If we choose δ such that

$$\frac{\epsilon}{4} - \frac{bb_2^2C(m,n)}{b_1^2\cos^4(\sqrt{b_1}R)}\delta^4 - \frac{\sqrt{b}b_3C(m,n)}{b_1\cos^2(\sqrt{b_1}R)}\delta^4 - b_2^2C(m,n)\delta^4 = 0,$$

then $\|\Psi\|^2$ is a bounded subharmonic function on M .

Now for any positive number c , let $u = (\|\Psi\|^2 + c)^{-1/2}$, then

$$\begin{aligned} \Delta u &= -\frac{1}{2}(\|\Psi\|^2 + c)^{-3/2}\Delta\|\Psi\|^2 + \frac{3}{4}(\|\Psi\|^2 + c)^{-5/2}\|\nabla\|\Psi\|^2\|^2 \\ (3.1) \quad &= -\frac{1}{2}(\|\Psi\|^2 + c)^{-3/2}\Delta\|\Psi\|^2 + 3(\|\Psi\|^2 + c)^{1/2}\|\nabla u\|^2 \\ &\leq -\frac{C_0}{2}(\|\Psi\|^2 + c)^{-3/2}\|\Psi\|^2 + 3(\|\Psi\|^2 + c)^{1/2}\|\nabla u\|^2. \end{aligned}$$

Since the Ricci curvature of M is nonnegative, the Omori-Yau maximum principle holds (cf. [7, 14]), that is, for any $\eta > 0$, there exists a point $p \in M$ such that at p ,

$$u < \inf u + \eta, \quad \|\nabla u\| < \eta, \quad \Delta u > -\eta.$$

It follows from $\|\Psi\| \leq \delta$ and (3.1) that

$$\frac{C_0}{2}\|\Psi\|^2 < \eta \left(\inf(\|\Psi\|^2 + c)^{-\frac{1}{2}} + 4\eta \right) (\|\Psi\|^2 + c)^2 < \eta (\delta^{-1} + 4\eta) (\delta^2 + c)^2.$$

Let $\eta \rightarrow 0$, we obtain $\sup\|\Psi\| = 0$.

Since Theorem 2.1, we have $\|d\Phi\| = 0$. Therefore Φ is a constant. \square

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